

***Formation of atomic hydrogen beam and intensity
limitations of atomic beam type polarized ion
sources and gaseous targets***

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Content:

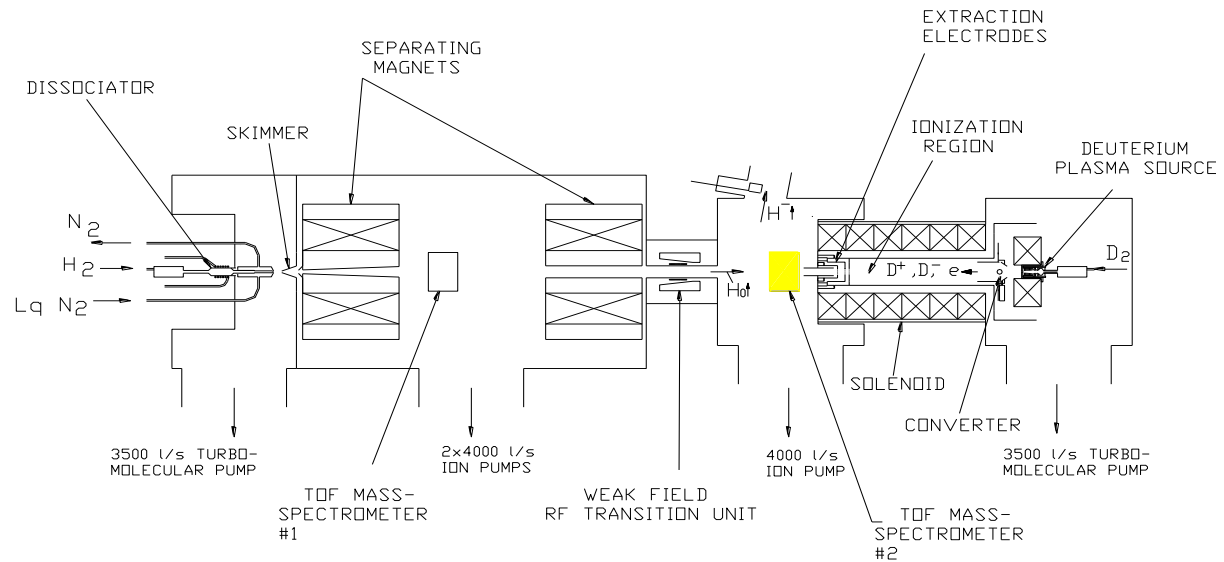
1. Intensity limitation of atomic hydrogen beams
2. Formation of atomic/molecular beams
3. Measurements of radial size of atomic hydrogen beam source at IKP, Juelich
4. Summary

Study of polarized pulsed atomic hydrogen beam at INR RAS

H- current -
4mA

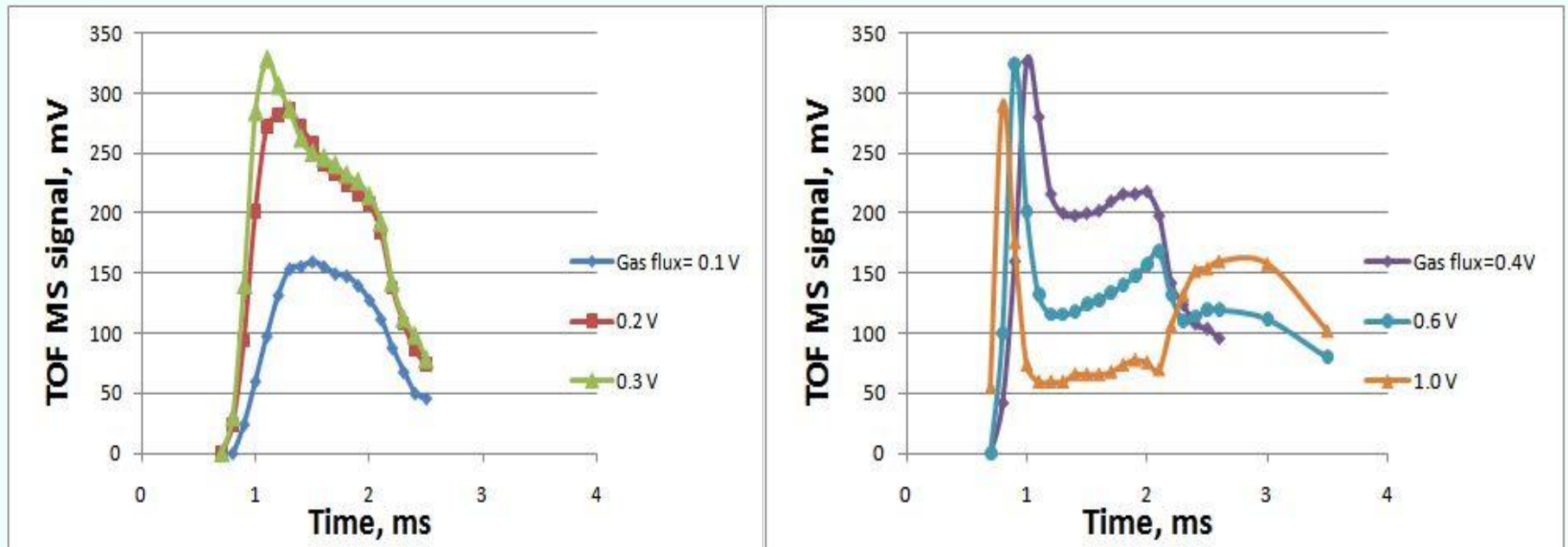
Polarization -
90%

Pulsed
operation,
200 μ s, 10 Hz

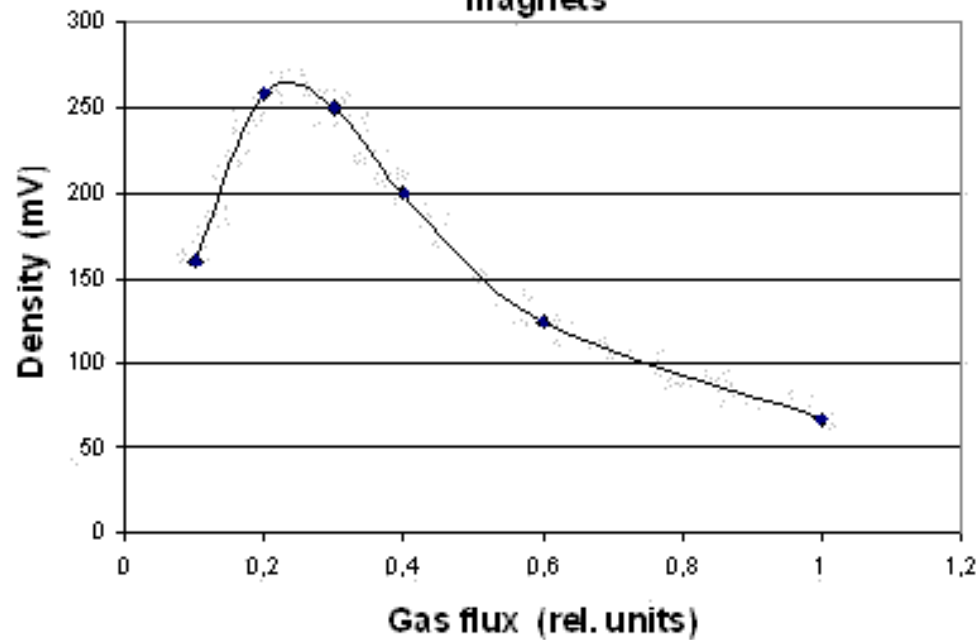


Schematic layout of the apparatus for study of pulsed polarized atomic hydrogen beam

Pulsed polarized atomic hydrogen beam density measurements (INR source)



Density of atomic hydrogen beam vs gas flux downstream the sextupole magnets

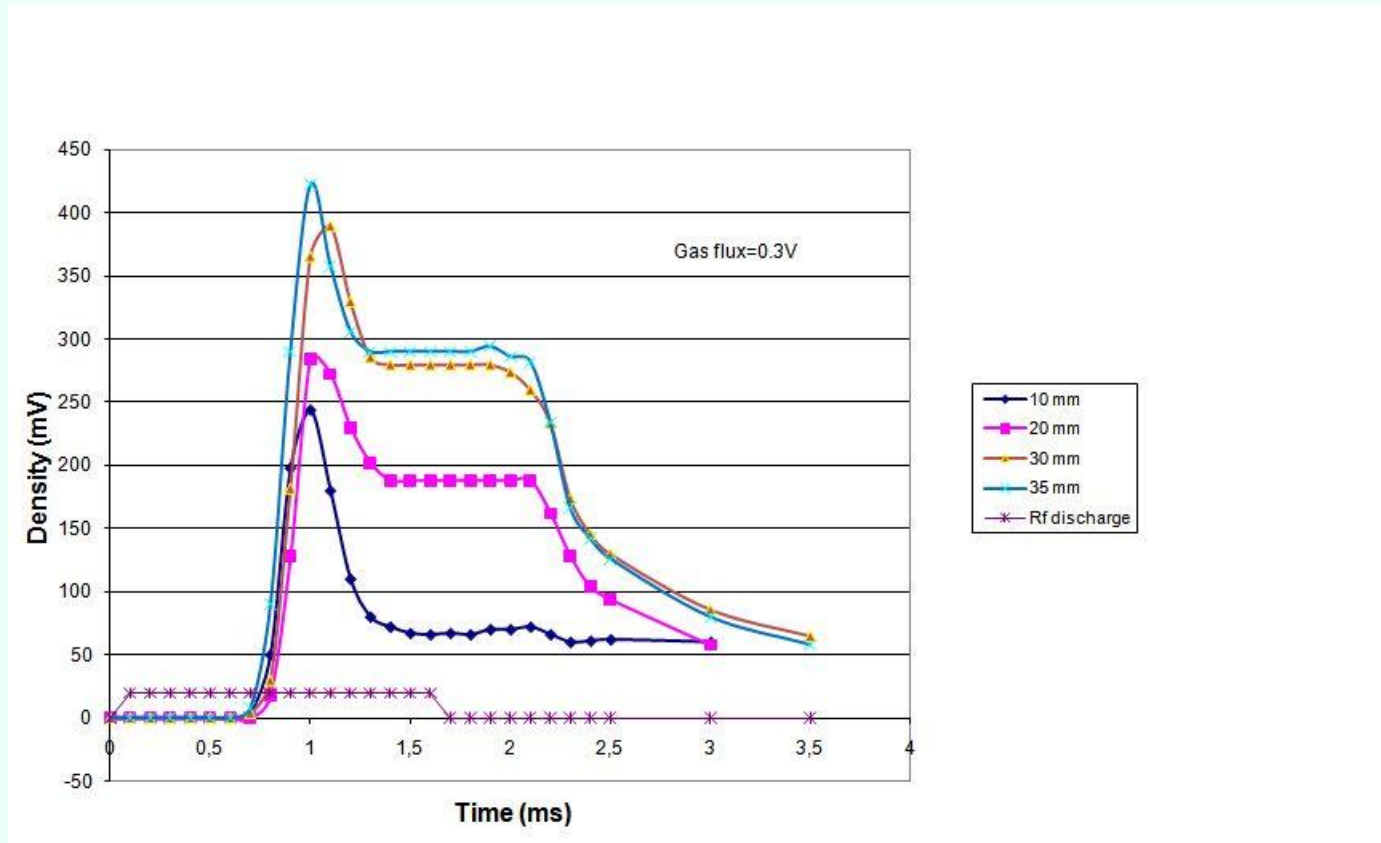


What is the reason for the pulsed atomic hydrogen beam intensity limitation?

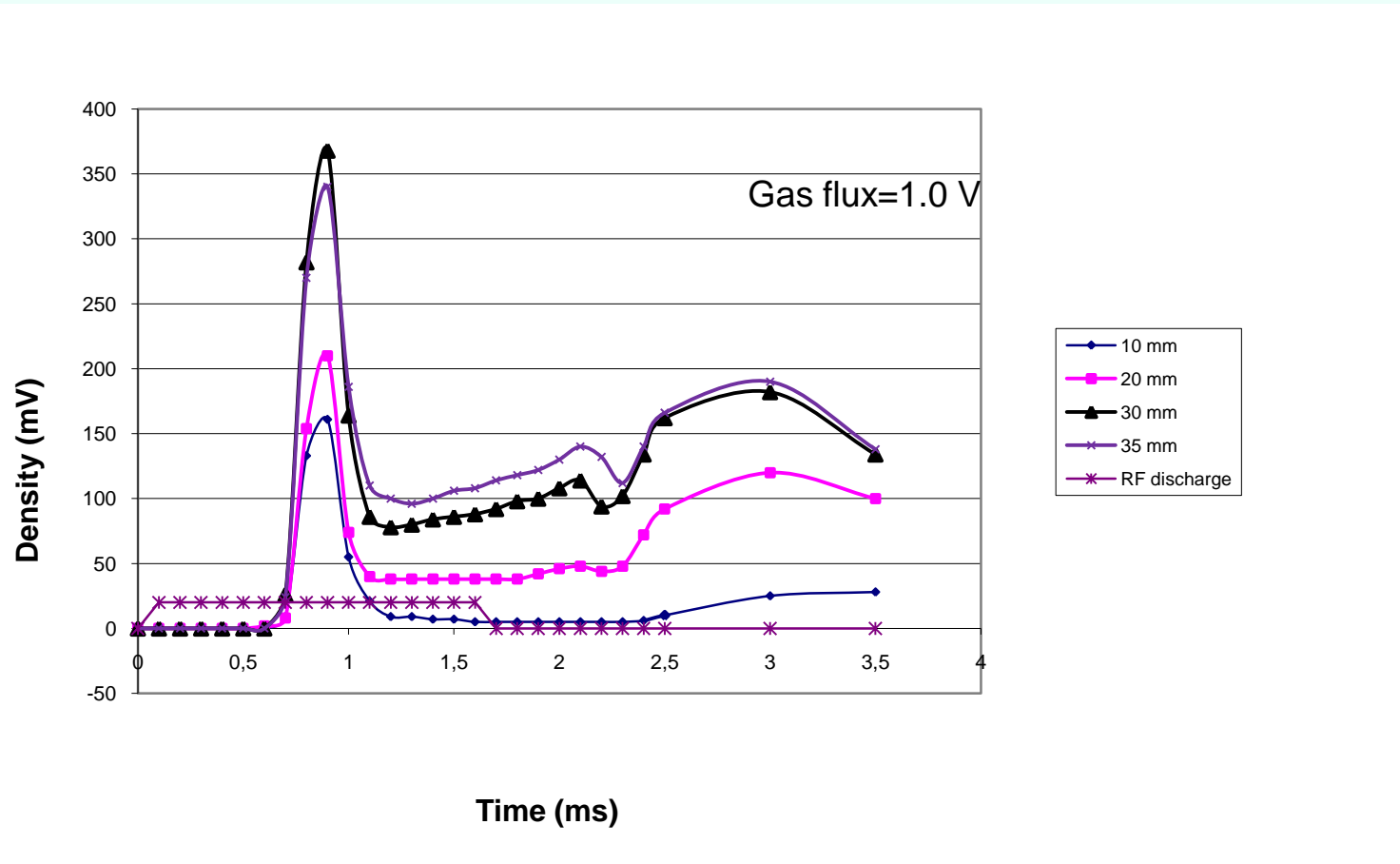
Parameters of velocity distribution change with change of gas flux:

- With increase of gas flux recombination increases
- This can be partially overcome by increase of RF discharge power
- But increase of RF power leads to velocities increase → focusing of faster atoms changes

**But we can fix parameters of the RF discharge and change only distance between the skimmer and the dissociator nozzle:
strong attenuation of atomic hydrogen beam takes place at small distances between the nozzle and the skimmer**



Atomic hydrogen beam dynamic for different distance between skimmer and dissociator nozzle (for big gas flux)

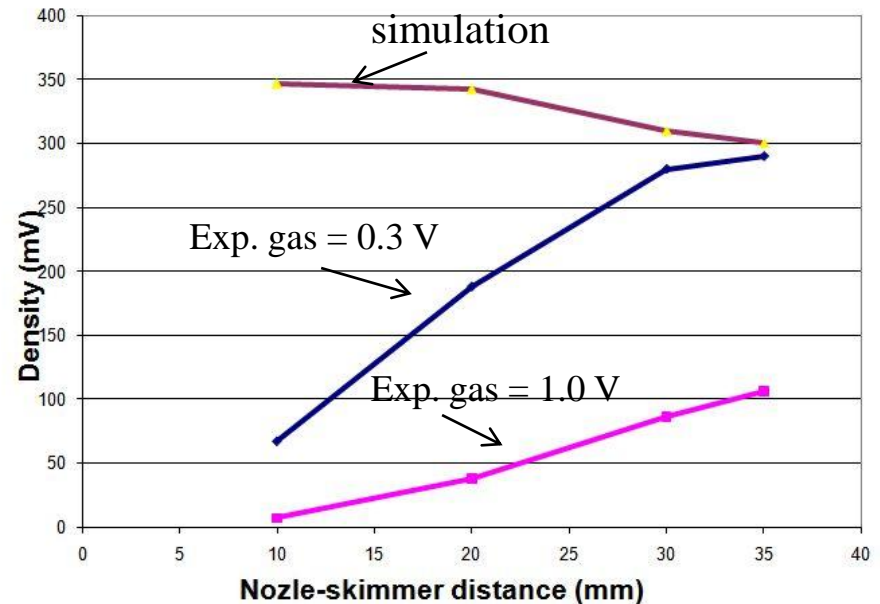


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Density of polarized atomic hydrogen beam (at middle of a pulse) VS nozzle-skimmer distance

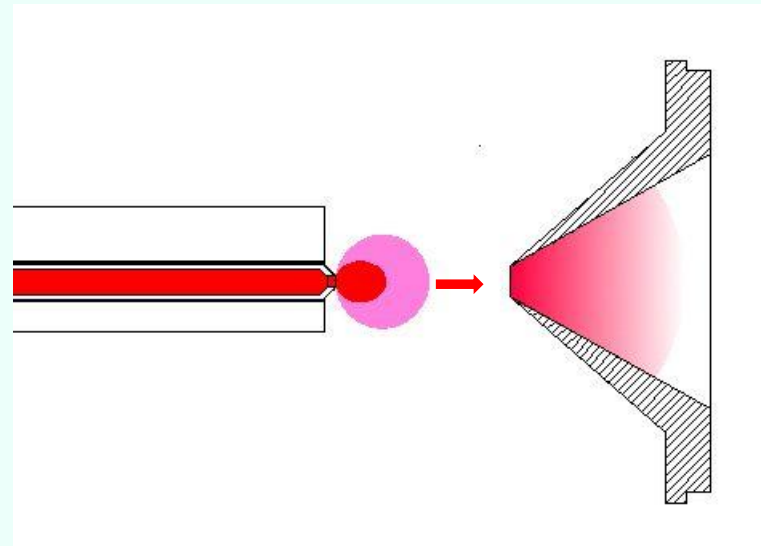
- Density of **pulsed** atomic hydrogen beam decreases with decrease of the nozzle-skimmer distance while simulation (which does not take into account scattering of atoms) forecasts increase of density



Possible explanation of the observed effects

(A. Belov –PSTP2007):

- Hydrogen gas “cloud” is formed inside the skimmer due to restricted skimmer conductance
- Atomic hydrogen beam is attenuated passing through the “cloud” due to scattering
- The “cloud” should produce secondary beam with size larger than the skimmer diameter



Formation of atomic/molecular beams

Effusion: $Kn \geq 1$ (mfp \geq orifice diameter)

- Cosine angle distribution
- Forward intensity:

$$I_0 = n_0 v A / 4\pi \quad (\text{sr}^{-1}\text{s}^{-1})$$

n_0 is the density of particles in the source
 v is the average velocity $v = (8kT_0/\pi m)^{1/2}$
 A is the orifice area

- Maximum intensity for effusive sources ($Kn \sim 1$)

$$I_{0\max} \sim 10^{18} \text{ s}^{-1} \text{ ster}^{-1}$$

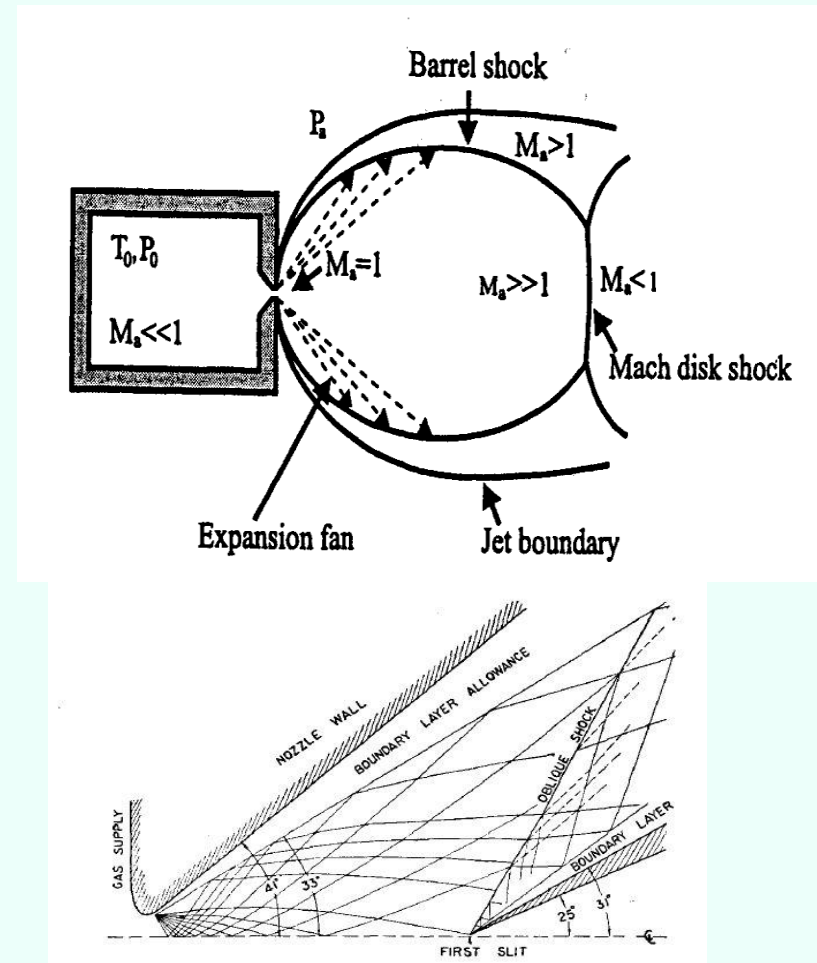
$$(T=80 \text{ K}, d_n = 2 \text{ mm}, n_{0\max} \sim 10^{15} \text{ cm}^{-3})$$

- For effusive sources:

source diameter = dissociator orifice diameter

Supersonic nozzle

- Use of supersonic nozzle for production of molecular beam has been proposed by A. Kantrowitz and J. Grey in 1951
 $Kn \ll 1$ ($mfp \ll$ nozzle diameter)
- Gaseous jet expansion through a sonic nozzle to a gas with ambient pressure P_a
- Efficient cooling during the jet expansion
- Complex phenomena including shock formation



Supersonic nozzle in ABS

- **Hydrogen flux** $\sim 1 \text{ mbar l s}^{-1}$ ($5 \cdot 10^{19} \text{ s}^{-1}$), $\text{Kn} \sim 0,02$
- **Density in an expanding hydrogen jet drops with distance from a nozzle as:**

$$n(z) \cong n_0 / (1+z^2/r_n^2)$$

Density decreases two orders of magnitude at distance $z/r_n \cong 10$. After this point collisions become rare (freezing surface) and jet becomes source of molecular beam

- **For supersonic nozzle source:**
 - Source diameter is determined by freezing surface geometry**
- **D. K. Toporkov (2004) suggested that freezing surface effective diameter can be larger than a skimmer diameter (screening effect)**

- **Intra jet scattering does not stop abruptly. It changes the source size too.**
- **For condition of ABS normal density shock (Mach disk shock) moves far away from a nozzle**
 $Z_M/d \sim \sqrt{P_0/P_a}$
for $P_0 \sim 1$ mbar, $P_a \sim 10^{-4}$ mbar
 $Z_M/d \sim 10^2$
- **DMCS does not show another density shocks formation in ABS**

Measurements of radial size of atomic hydrogen beam

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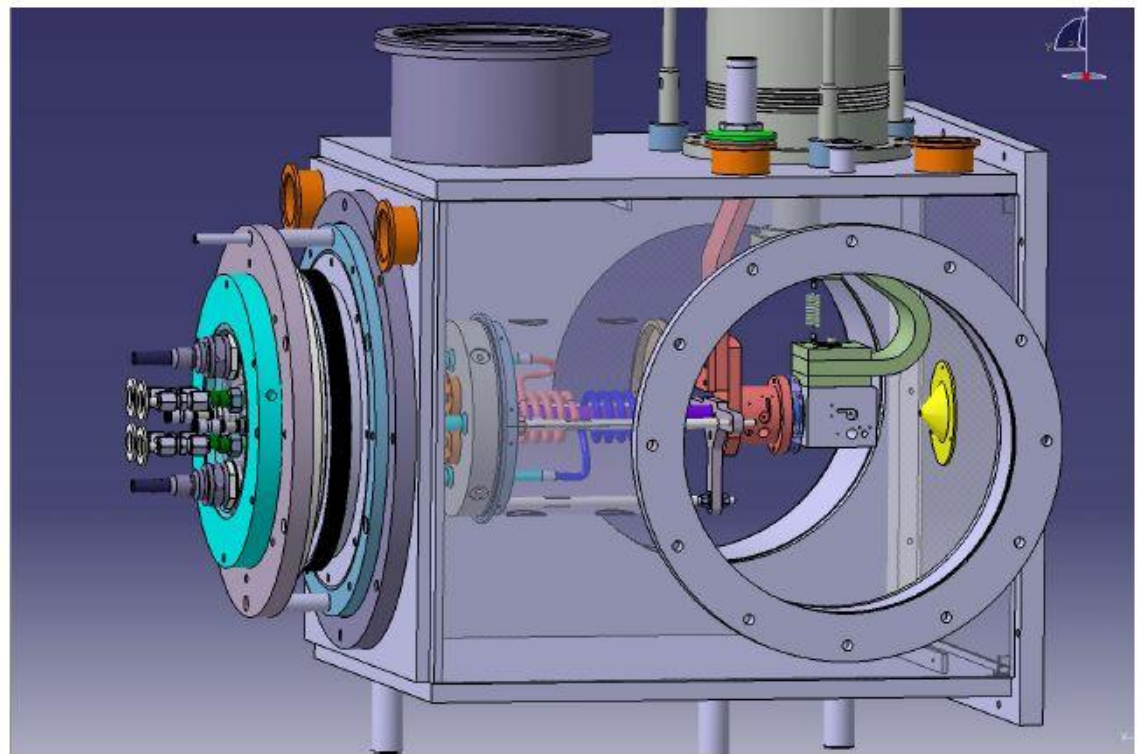
***INR RAS, Moscow, Russia**

****IKP, Juelich, Germany**

COSY – Juelich pulsed RF discharge dissociator

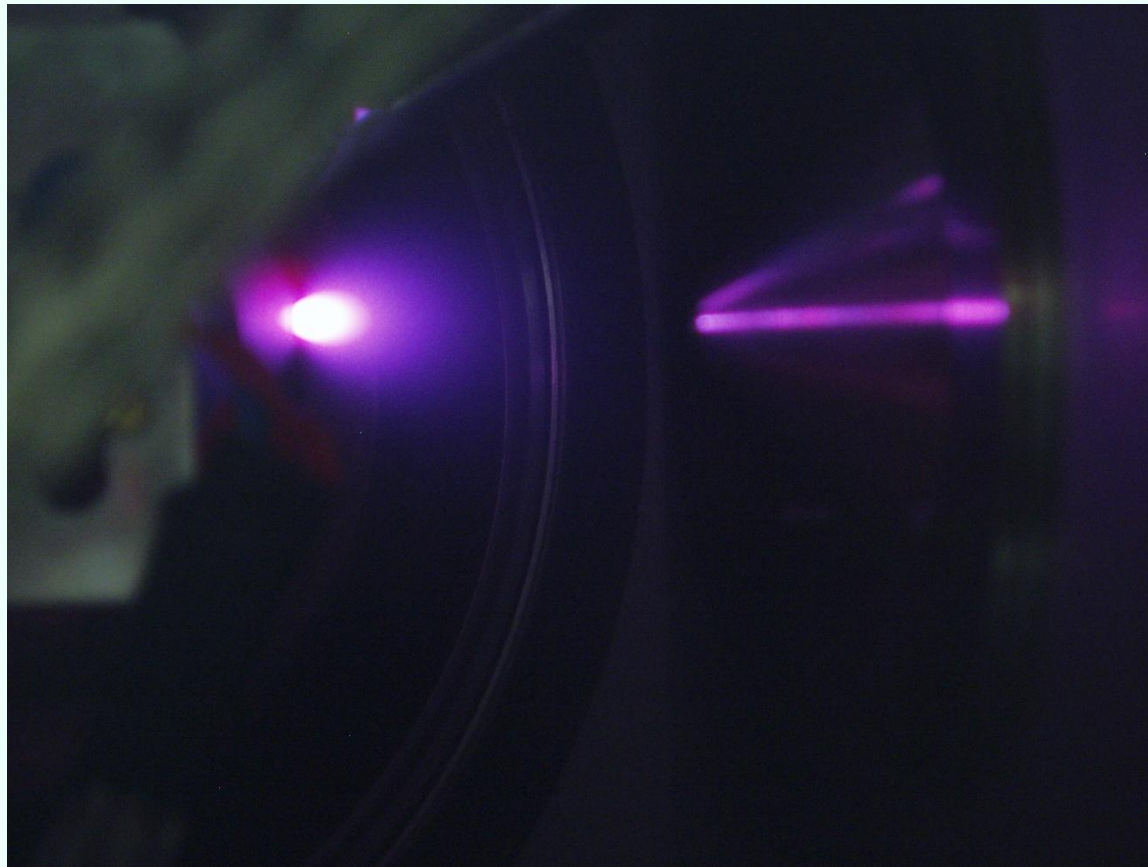
RF pulse duration
~20 ms

Rep. rate – 1 Hz



Nozzle radius =1.25 mm, Skimmer radius =3 mm,
nozzle-skimmer distance =35 mm

Photo picture of nozzle-skimmer region of COSY test-bench dissociator

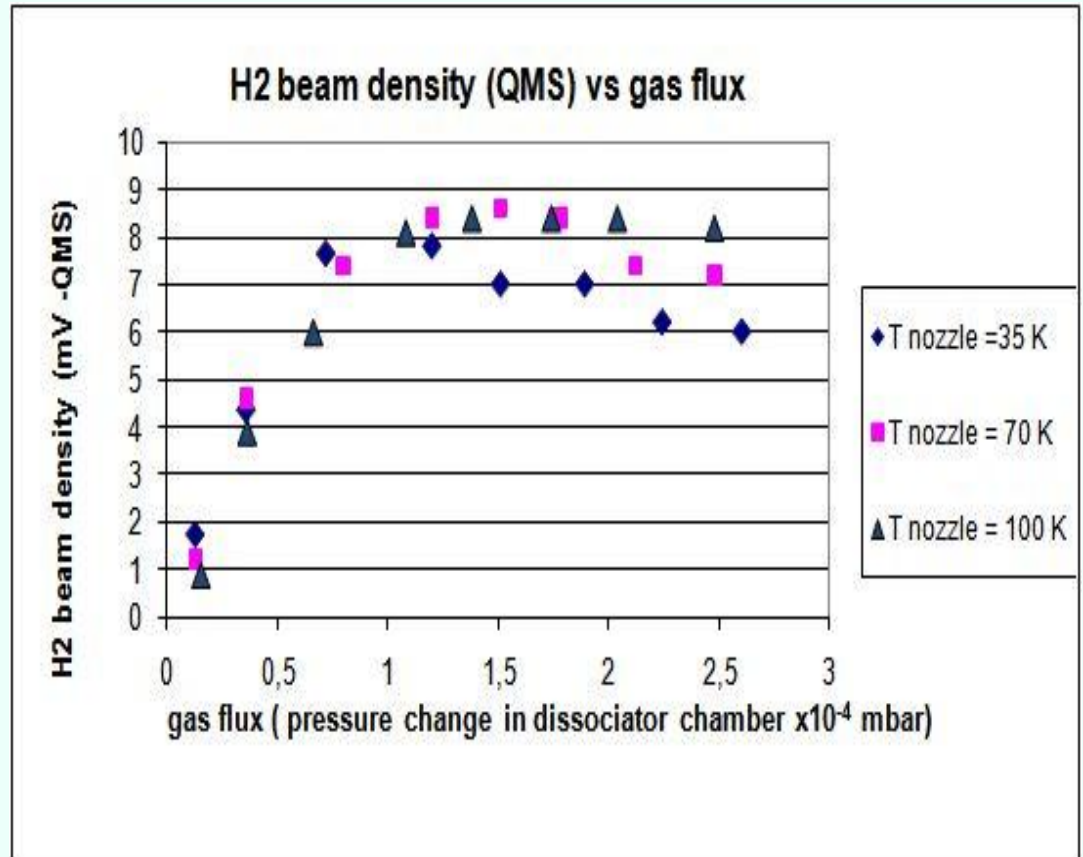


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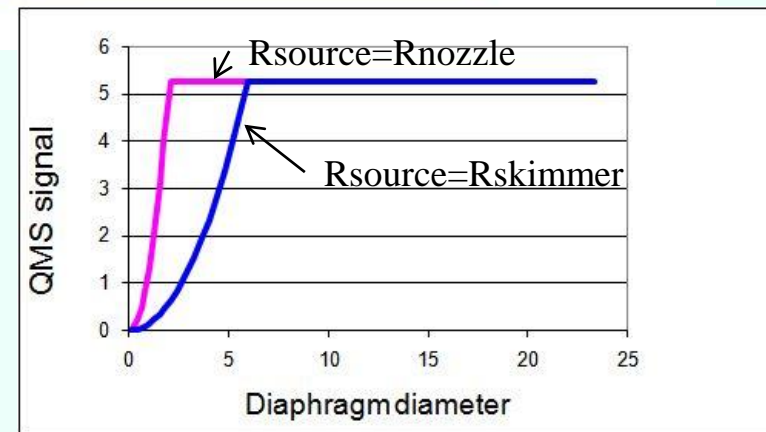
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- Molecular hydrogen beam density is restricted in spite of pulsed mode of gas supply

- H₂ beam density saturation takes place at gas flux (pressure change in dissociator vacuum chamber) = (1-1.5) 10⁻⁴ mbar

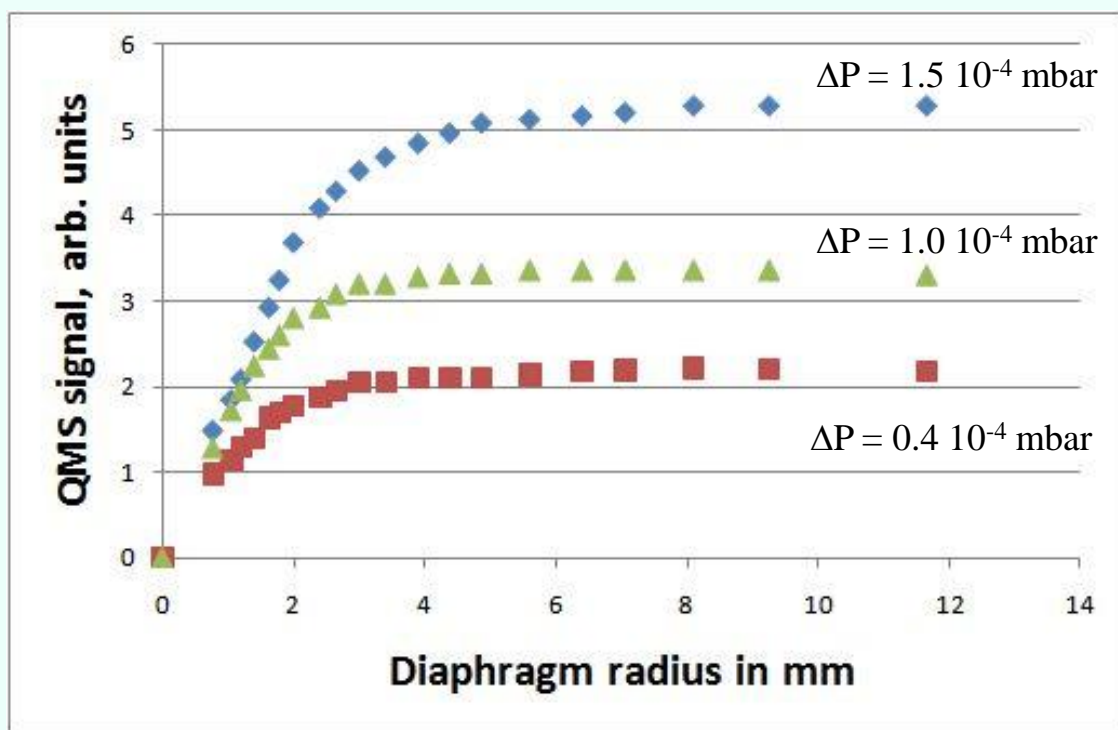


Schematic of apparatus for hydrogen beam source size measurements



Results of measurements of atomic hydrogen source size for different molecular hydrogen gas flux with diaphragm method

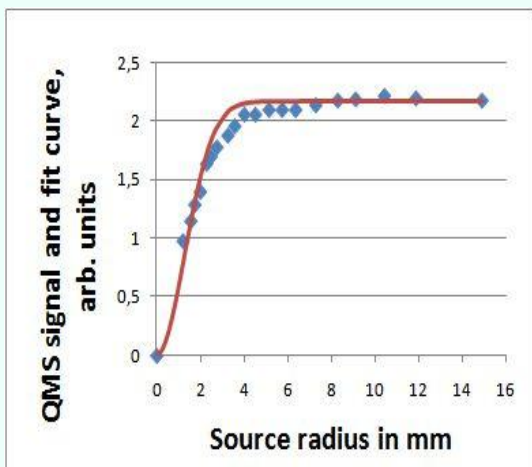
Nozzle radius =1.25 mm, Skimmer radius =3 mm, H atoms



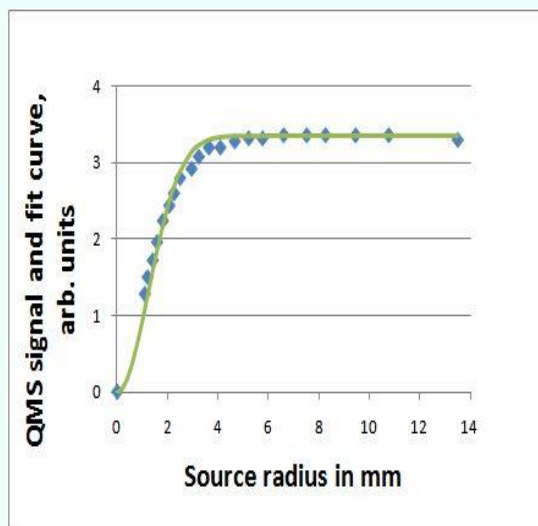
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Fit of data by Gaussian distribution

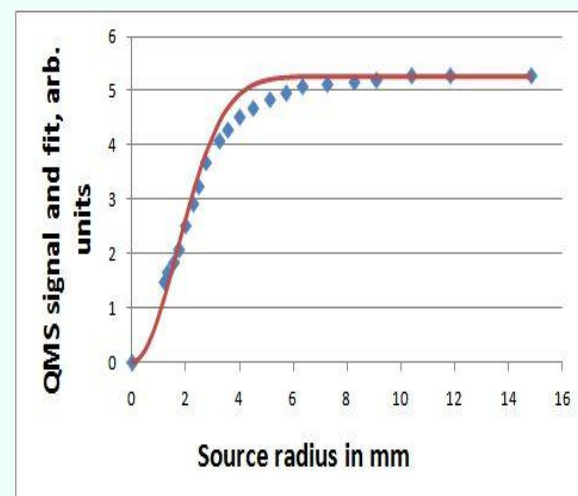
Nozzle radius =1.25 mm, Skimmer radius =3 mm, H atoms



$\Delta P = 0.4 \cdot 10^{-4}$ mbar,
R=1,74 mm



$\Delta P = 1.0 \cdot 10^{-4}$ mbar,
R=1,78 mm



$\Delta P = 1.5 \cdot 10^{-4}$ mbar,
R=2,4 mm

Nozzle radius =1.25 mm,
Skimmer radius =3 mm,
H atoms

Comparison of

Fit of data by one gaussian
distribution:

$$F = A_1 (1 - \exp(-r^2/R_1^2))$$

$$R_1 = 2.4 \text{ mm}$$

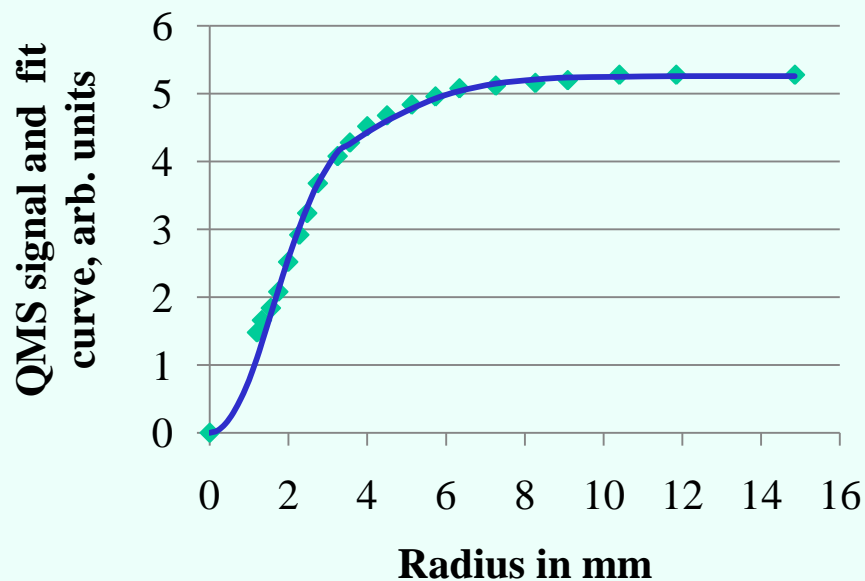
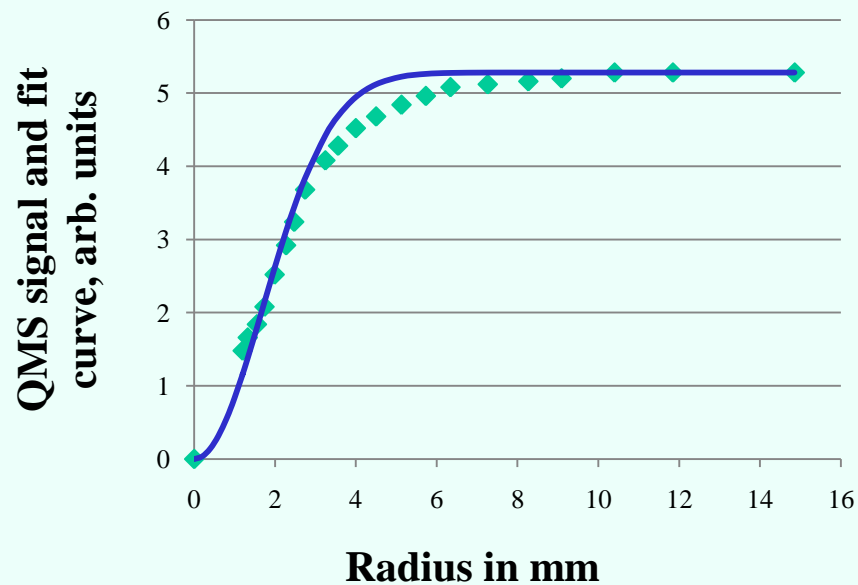
and

Fit of data by two gaussian
distributions:

$$F = A_1 (1 - \exp(-r^2/R_1^2)) + \\ + A_2 (1 - \exp(-r^2/R_2^2))$$

$$R_1 = 1.9 \text{ mm} \quad R_2 = 4.22 \text{ mm}$$

$$A_1 = 3.6 \quad A_2 = 1.94$$



Nozzle radius =1.25 mm,
Skimmer radius =3 mm,
H atoms

Comparison of

Fit of data by one gaussian
distribution:

$$F = A_1 (1 - \exp(-r^2/R_1^2))$$

$$R_1 = 2.4 \text{ mm}$$

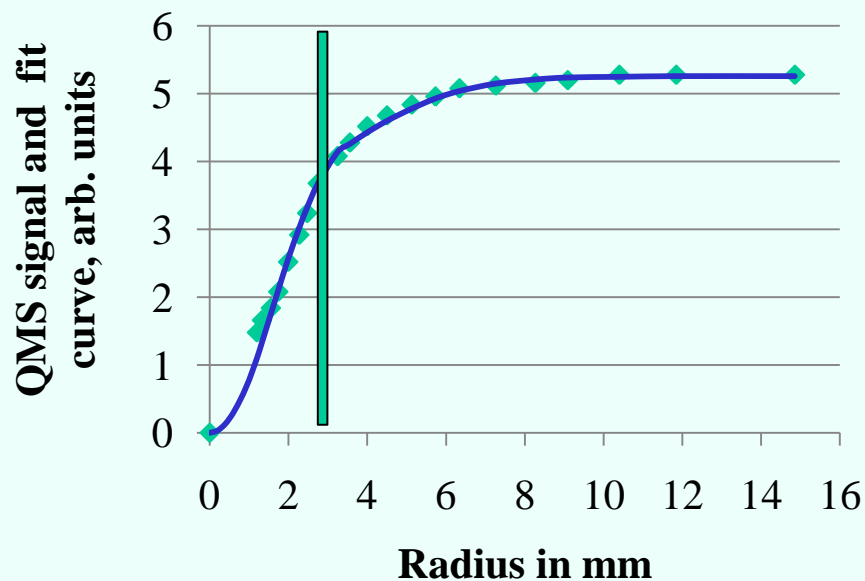
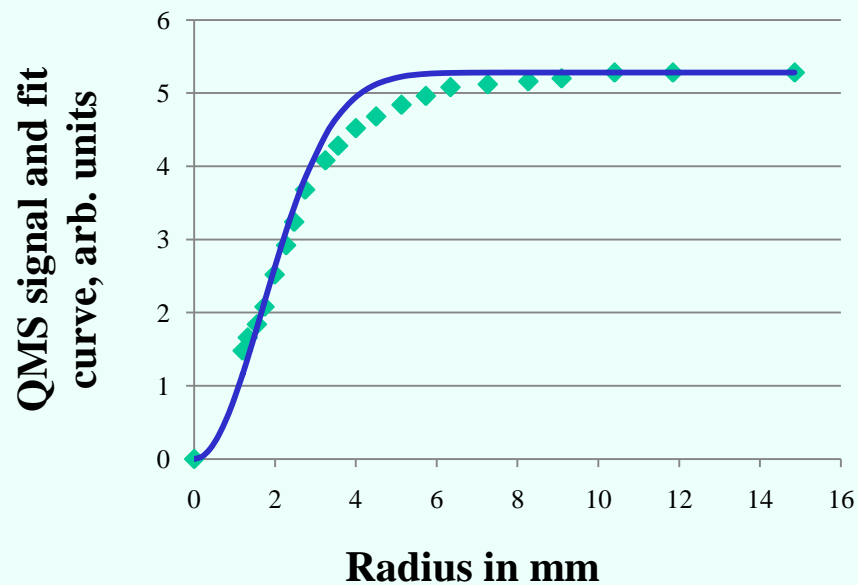
and

Fit of data by two gaussian
distributions:

$$F = A_1 (1 - \exp(-r^2/R_1^2)) + \\ + A_2 (1 - \exp(-r^2/R_2^2))$$

$$R_1 = 1.9 \text{ mm} \quad R_2 = 4.22 \text{ mm}$$

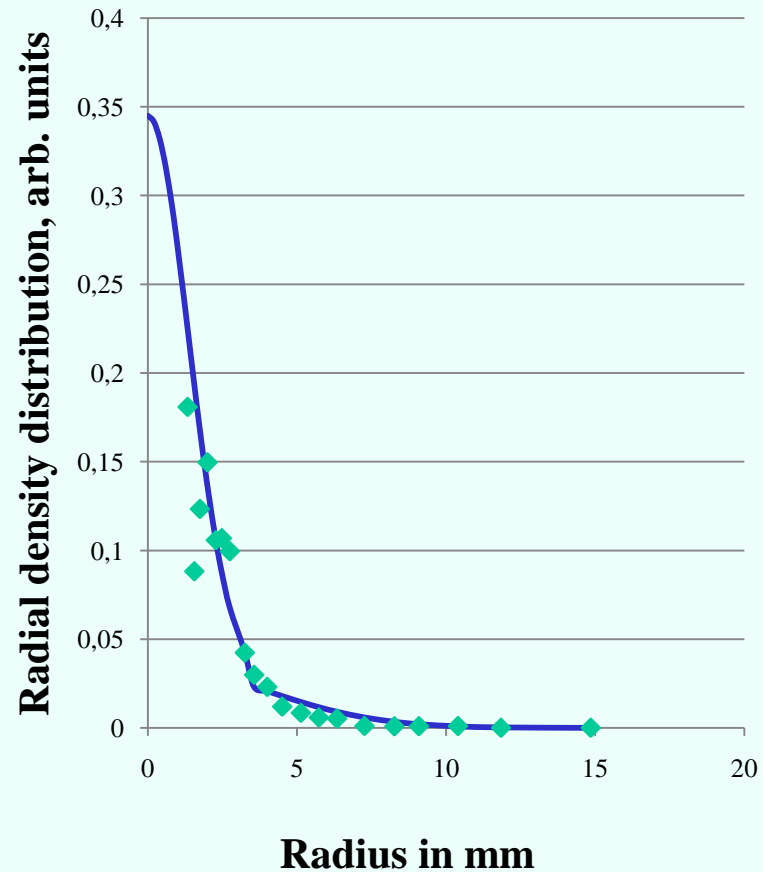
$$A_1 = 3.6 \quad A_2 = 1.94$$



Nozzle radius =1.25 mm, Skimmer radius =3 mm, H atoms

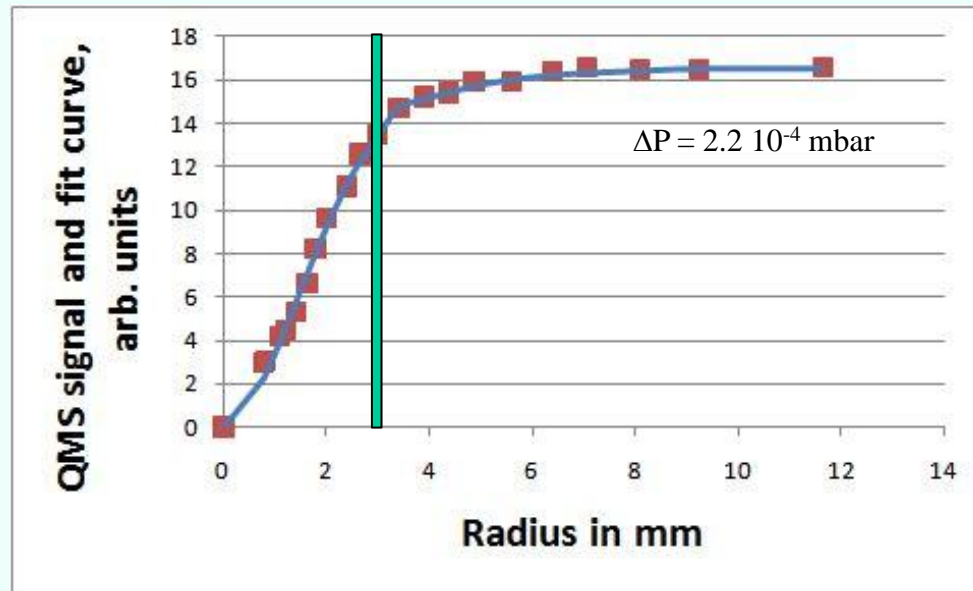
Radial density distribution
(fit by two gauss curve):

$$n(r) = A_1 \exp(-r^2/R_1^2) + A_2 \exp(-r^2/R_2^2)$$



Results for hydrogen molecular beam

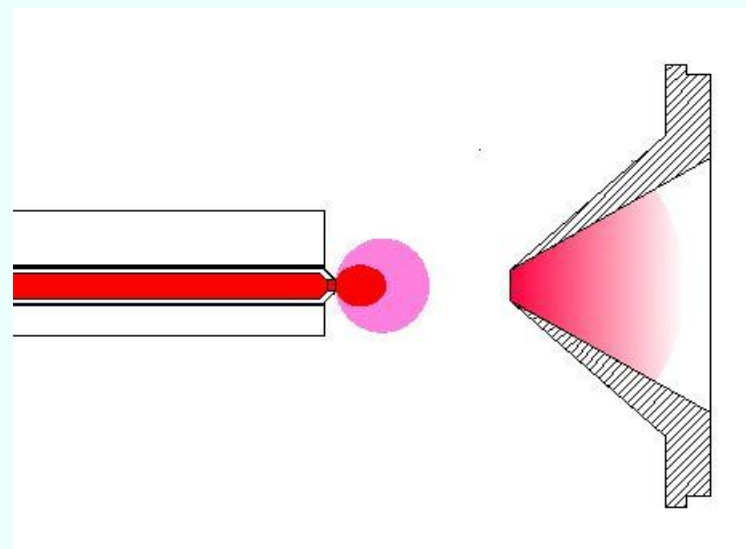
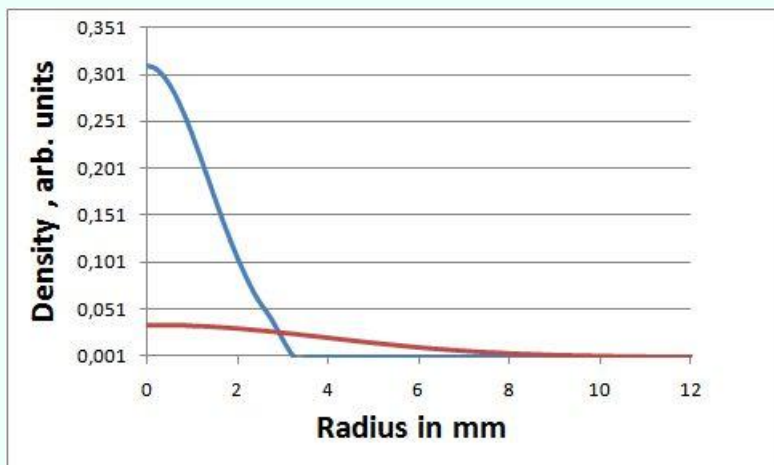
- Nozzle radius = 1.25 mm, Skimmer radius = 3 mm, H₂ molecules



$R_1 = 2.4$ mm, $R_2 = 4.5$ mm

The measurements confirm suggestion concerning attenuation of atomic hydrogen beam from ABS due to scattering on gas cloud inside the skimmer:

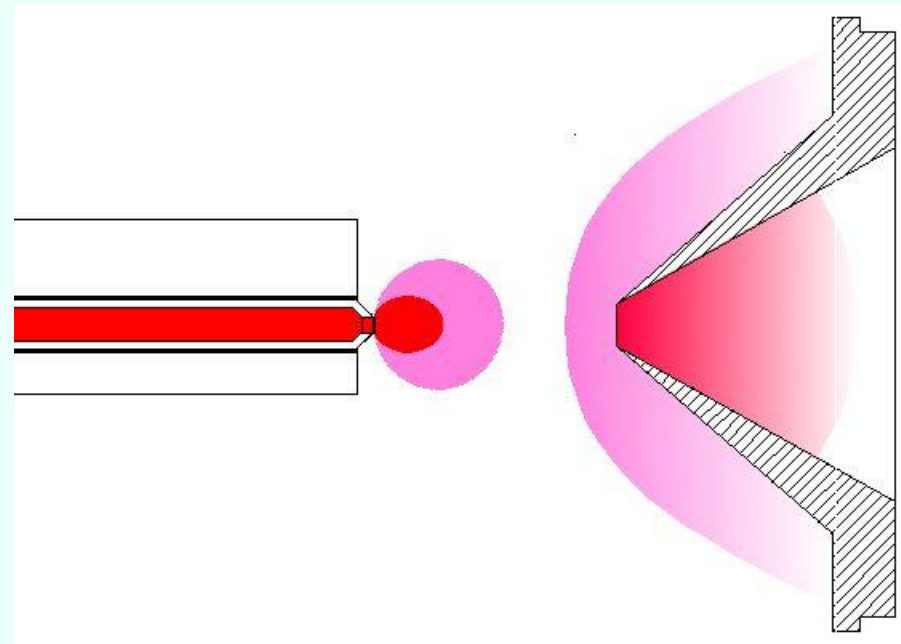
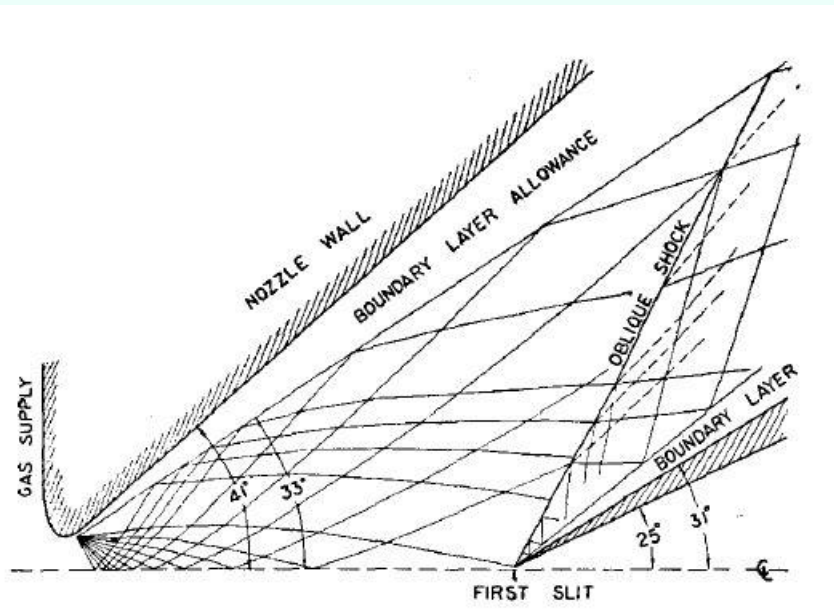
There are two sources of hydrogen atoms, one upstream the skimmer and another downstream the skimmer.



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- **Secondary source arises probably due to density shock at the skimmer orifice which is produced by an intense hydrogen jet**



How to avoid the density shock formation?

- Increase nozzle-skimmer distance (for INR source this works well)

However,

- sextupole magnet geometry should be optimized for larger distance
- scattering on residual gas in dissociator chamber will increase.

- Optimize skimmer geometry

However,

- It is important to see the effect in simulation. Up to now DMCS does not show formation of density shocks.

Summary

- Characteristic radial size of atomic hydrogen beam source of COSY ABS increases with gas flux from 1.7 mm to 1.9 mm (while nozzle radius was 1.25 mm and skimmer radius – 3 mm).
- Tails in radial distribution of the atomic hydrogen source density with size larger than skimmer radius appear at high gas flux. This is connected probably with formation of density shock around skimmer orifice and gas cloud inside the skimmer. By this way secondary source of hydrogen atoms arises with radial size larger than the skimmer radius.
- Limitation of intensity of a ABS can be connected with scattering of atomic hydrogen beam on a density shock at a skimmer.
- It is important to take into account radial size of atoms source for optimization of ABS apparatus .